

Communications to the Editor

Exchangeable Proton NMR without Base-Line Distortion, Using New Strong-Pulse Sequences

Pierre Plateau and Maurice Guéron*

Groupe de Biophysique, Ecole Polytechnique
91128 Palaiseau Cedex, France

Received August 5, 1982

Pulsed nuclear magnetic resonance of exchangeable protons in H₂O solvent is most often accomplished by using Redfield's "2-1-4" weak-pulse sequence, which leaves the water proton magnetization unperturbed and provides second-order cancellation: the amplitude of a peak is a quadratic function of its distance to the frequency of cancellation.^{1,2} Hence, solvent suppression is insensitive to slight errors and drifts in frequency or field-gradient correction. The same features can be achieved by a sequence³ of strong, nonselective pulses (15°, τ , 15°, 15°, τ , 15°). The advantages of strong-pulse irradiation lie in the easier designing of sequences, in their insensitivity to long-term drift in pulse amplitude, and in the simpler generation of such pulses in an FT spectrometer.

A defect of the sequences mentioned above is the generation of phase shifts, which combined with the existence of a residual but still intense solvent peak, lead to base-line undulation.³ We have searched for strong-pulse solvent-suppressing sequences that would produce no phase shift and hence no base-line undulation. Two sequences will be presented here. They share the following properties: (a) only strong, nonselective pulses (10 W) are used; (b) the pulse is at the solvent frequency; (c) they use only two phases, 0° and 180°; (d) no linear phase correction is required. The first sequence ("jump and return") provides first-order cancellation of the solvent peak. The phase is strictly constant across the spectrum, except for a step change of 180° at the solvent frequency. The second sequence ("set, jump, and return") gives second-order cancellation and only small phase shifts.

The JR ("Jump and Return") Sequence (90°_y, τ , 90°_{-y}). The first pulse brings all spins along 0x in the rotating frame. During the waiting time they fan out in the xy plane; the solvent spins are at resonance and remain along 0x. The waiting time determines the variation of amplitude with frequency and the final longitudinal magnetization. The second pulse brings all spins back from the x0y plane to positions in the z0y plane, the solvent being returned to 0z. The free precession is then recorded. Transverse magnetizations start along 0y on one side of the solvent peak, along -0y on the other.

Tuning is done while observing the free precession. It consists in (a) adjusting the pulse length to 90° (an error will only give small phase shifts, it will not affect solvent cancellation), (b) correcting for pulse-phase errors by a slight change in radiofrequency (rf) or by an adjustable phase differential, (c) correcting for pulse inequality (and/or radiation damping) by an adjustable pulse length differential.

Figure 1 shows a spectrum of yeast tRNA^{Phe} obtained with the JR sequence. It exhibits the characteristic 180° phase difference between the two sides of the solvent peak. Good phase control and shaping of the rf pulses are involved in these experiments. In our spectrometer, all rf manipulations are carried out at the

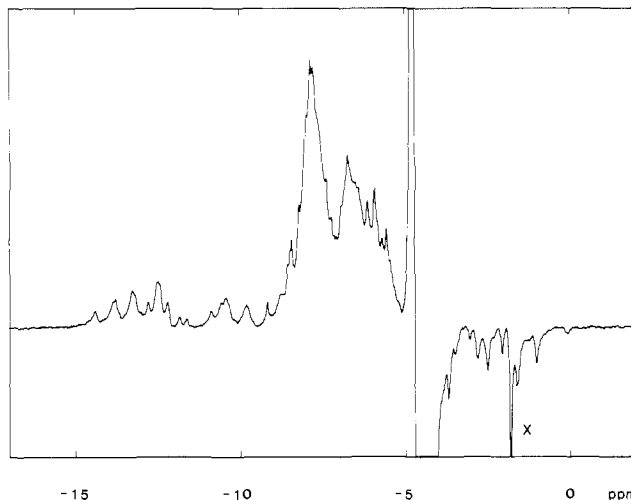


Figure 1. 276-MHz JR spectrum of yeast tRNA^{Phe}. Conditions: 0.9 mM tRNA; 35 °C; 150 FIDs; waiting time τ between the two pulses of the JR sequence, 125 μ s; repetition time, 0.2 s; broadening, 5 Hz. No phase correction has been applied except for audio-filter response. The water peak (at -4.65 ppm) is reduced 500-fold. X is the cacodylate buffer. Note the good signal-to-noise ratio.

intermediate frequency of 10 MHz.⁴

Water is reduced by a factor of 300 on each free precession, similar to what we obtain with 2-1-4 type sequences. It is apparently limited by phase fluctuations of the rf pulses. Phase hum or fluctuations of the generator show up clearly in the free precession of a strong signal, at resonance, in the dispersion mode. They give deviations from zero. Hum in the field shimming coil shows up similarly.

The SJR ("Set, Jump, and Return") Sequence (θ°_{-y} , 10τ , θ°_{y} , JR). The JR sequence provides first-order cancellation only. Sequences providing second-order cancellation without phase-shift are unknown to us; they may be restricted by mathematical considerations. The SJR sequence gives second-order cancellation at the cost of modest phase shifts.

We start with a pulse of tilt-angle $\theta = 5.7^\circ$ (i.e., $\sin \theta = 0.1$) along -0y and wait for 10 times longer than the JR waiting time. During this time, all spins precess on a small circle of axis 0z. Spins that are at resonance do not move. Those that are near enough for a linear approximation to be valid spread out along the tangent (parallel to the 0y axis) by the same amount as they will during the waiting time of the upcoming JR sequence, *but in the opposite sense*. Hence they will be refocused by JR.

The range of second-order cancellation is changed by varying the "set" angle θ , the first waiting period being set to $(1/\sin \theta)$ times the JR one. Spins further away in frequency precess all around the small circle. This will modify their position at the end of the sequence, inducing phase shifts. These are large close to the solvent but decrease and oscillate as one goes away from the solvent frequency. Figure 2 compares the JR and SJR spectra.

SJR has been used for various NOE experiments in which selective line inversion is also obtained by the use of strong pulses (in preparation). Pulse sequences devoid of phase-shifts could be particularly useful in 2-D NMR of exchangeable protons. Respect of spectral amplitudes is also important in 2-D NMR. The sine amplitude response of the JR sequence (and SJR as well,

(1) Redfield, A. G.; Kunz, S. D.; Ralph, E. K. *J. Magn. Reson.* **1975**, *19*, 114-117.

(2) Redfield, A. G.; Kunz, S. D. "NMR in Biochemistry"; Opella, S. J., Lu, P., Eds.; Marcel Dekker: New York, 1979; pp 225-239.

(3) Plateau, P.; Dumas, C.; Guéron, M., manuscript in preparation.

(4) Caron, F.; Guéron, M.; Nguyen Ngoc Quoc Thuy; Herzog, R. F.; *Rev. Phys. Appl.* **1980**, *15*, 1267-1274.

